## ECOLE POLYTECHNIQUE FEDERALE DE LAUSANNE

School of Computer and Communication Sciences

Handout 3	Statistical Physics for Communication and Computer Science
Homework 3	March 10, 2011, INR 113 - 9:15-11:00

Last time you proved that the Ising model in one dimension (d = 1) does not have a phase transition for any T > 0. On the grid  $\mathbb{Z}^d$  there is a non trivial phase diagram with first and second order phase transitions for any  $d \ge 2$ . This is also the case on the complete graph (as shown in the lectures) which morally corresponds to  $d = +\infty$ . Another graph that in a sense, corresponds to  $d = +\infty$ , is the q-ary tree. Indeed on  $\mathbb{Z}^d$  the number of lattice sites at distance less than n from the origin scales as  $n^d$ . On the q-ary tree it scales as  $(q-1)^n$  which grows faster than  $n^d$  for any finite d.

The goal of the two exercises below is to solve for the Ising model on a q-ary tree and show that it displays first and second order phase transitions (with similar qualitative properties than on a complete graph).

Consider a finite rooted tree and call the root vertex o. All vertices have degree q, except for the leaf nodes that have degree 1. We suppose that the tree has n levels (the root being "level 0"). The thermodynamic limit corresponds to  $n \to +\infty$ . The Hamiltonian (multiplied by  $\beta$ ) is

$$\beta \mathcal{H}_n = -K \sum_{(i,j)\in E_n} s_i s_j - h \sum_{i\in V_n} s_i \tag{1}$$

were K > 0,  $h \in \mathbb{R}$ ,  $V_n$  is the set of vertices and  $E_n$  the set of edges. We are interested in the magnetization of the root node in the thermodynamic limit:

$$m(K,h) = \lim_{n \to +\infty} \langle s_o \rangle_n = \frac{\sum_{\{s_k, k \in V_n\}} s_o e^{-\beta \mathcal{H}_n}}{Z_n}$$
(2)

The formula  $\tanh^{-1} y = \frac{1}{2} \ln \frac{1+y}{1-y}$  might be useful.

**Problem 1** (Recursive equations). Perform the sums over the spins attached at the leaf nodes and show that

$$\langle s_o \rangle_n = \frac{\sum_{\{s_k, k \in V_{n-1}\}} s_o e^{-\beta \mathcal{H}'_{n-1}}}{Z'_{n-1}}$$
(3)

where  $E_{n-1}$  and  $V_{n-1}$  are the edge and vertex sets of a tree with with n-1 levels and the new Hamiltonian is

$$\beta \mathcal{H}'_n = -K \sum_{(i,j) \in E_{n-1}} s_i s_j - h \sum_{i \in V_{n-1}} s_i - (q-1) \tanh^{-1}(\tanh K \tanh h) \sum_{i \in \text{level } n-1} s_i \qquad (4)$$

Iterate this calculation and deduce

$$\langle s_o \rangle_n = \tanh(h + q \tanh^{-1}(\tanh K \tanh u_n))$$
 (5)

where

$$u_{k+1} = h + (q-1) \tanh^{-1}(\tanh K \tanh u_k), \qquad u_1 = h$$
 (6)

Check that for q = 2 you get back the recursion of homework 2.

**Problem 2** (Analysis of the recursion). We want to analyze the fixed point equation for  $q \geq 3$ ,

$$u = h + (q - 1) \tanh^{-1}(\tanh K \tanh u) \tag{7}$$

Plot the curves  $u \to u - h$  and  $u \to (q - 1) \tanh^{-1}(\tanh K \tanh u)$  and show that:

- for  $K \leq K_c \equiv \frac{1}{2} \ln \frac{q}{q-2} = \tanh^{-1}(q-1)^{-1}$ , (7) has a unique solution, and that the iterations (6) converge to this unique solution.
- for  $K > K_c$ :
  - for  $|h| \ge h_s$ , (7) has a unique solution (you do not needw3 to compute  $h_s$  explicitly although it is possible to find its analytical expression) and that the iterations (6) converge to this unique solution.
  - for  $|h| < h_s$ , (7) has three solutions  $u_-(h) < u_0(h) < u_+(h)$ . Check graphically that for h > 0 the iterations (6) with initial condition  $u_1 = h$  converge to  $u_+(h)$ . Similarly for h < 0 they converge to  $u_-(h)$ . Check also graphically that the fixed point  $u_0(h)$  is unstable whereas  $u_{\pm}(h)$  are stable.

**Problem 3** (Phase transitions). Now we want to discuss the consequences of the results in problem 2 for the phase diagram. In a nutshell: in the  $(K^{-1}, h)$  plane there is a first order phase transition line  $(K^{-1} \in [0, K_c^{-1}], h = 0)$  terminated by a critical point  $K_c$ . Outside of this line m(K, h) is an analytic function of each variable.

We define the "spontaneous magnetization" as  $m_{\pm}(K) = \lim_{h \to 0_{\pm}} m(K, h)$ .

- Deduce from the analysis in problem 2 that for  $K \leq K_c$ ,  $m_+(K) = m_-(K) = 0$ .
- Deduce that for  $K > K_c$ ,  $m_+(K) \neq m_-(K)$  (jump discontinuity or first order phase transition) and that for  $K \to +\infty$   $m_{\pm} \to \pm 1$ .
- Show that for  $K \to K_c$  from above,  $m_{\pm}(K) \sim (K K_c)^{1/2}$ . So on the line h = 0, as a function of K, the spontaneous magnetization is continuous but not differentiable at  $K_c$  (second order phase transition).
- Now fix  $K = K_c$  and show that  $m(K_c, h) \sim |h|^{1/3}$ . As a function of h the spontaneous magnetization is continuous but not differentiable at  $K_c$  (second order phase transition).

**Hint:** for the last two questions you can expand the fixed point equation to order  $u^3$ .

**Remark:** Note that the exponents 1/2 and 1/3 are the same than for the model on a complete graph. This is also the case for all  $d \ge 4$  and is not the case for d = 2, 3.