Exercice 1 Heisenberg Uncertainty Principle.

1. Let $|\phi_{\lambda}\rangle = (A + i\lambda B) |\psi\rangle$. For $\lambda \in \mathbb{R}$, let us define $f(\lambda) = \langle \psi_{\lambda} | \psi_{\lambda} \rangle$. It is clear that for any $\lambda \in \mathbb{R}$, $f(\lambda) \ge 0$. We also have

$$f(\lambda) = \langle \psi | (A^{\dagger} - i\lambda^* B^{\dagger})(A + i\lambda B) | \psi \rangle = \langle \psi | (A - i\lambda B)(A + i\lambda B) | \psi \rangle$$

= $\langle \psi | A^2 | \psi \rangle + \lambda^2 \langle \psi | B^2 | \psi \rangle + i\lambda \langle \psi | (AB - BA) | \psi \rangle$
= $\langle \psi | A^2 | \psi \rangle + \lambda^2 \langle \psi | B^2 | \psi \rangle + \lambda \langle \psi | i[A, B] | \psi \rangle$,

where we used the Hermitian property of A and B and the fact that $\lambda \in \mathbb{R}$, thus $\lambda = \lambda^*$. First one can simply check that the operator i[A, B] is a Hermitian operator so the last term $\langle \psi | i[A, B] | \psi \rangle$ is real-valued so $f(\lambda)$ is real-valued (this was already clear). We see that $f(\lambda)$ is a second order polynomial in $\lambda \in \mathbb{R}$. As it is non-negative for every value of λ , it results that its discriminant must be negative or zero (for a quadratic equation $ax^2 + bx + c = 0$ the discriminant is defined by $\Delta = b^2 - 4ac$). Hence, we get

$$|\langle \psi | [A, B] | \psi \rangle|^2 \le 4 \langle \psi | A^2 | \psi \rangle \langle \psi | B^2 | \psi \rangle.$$

As we assumed that the operators have zero mean, $\langle \psi | A | \psi \rangle = \langle \psi | B | \psi \rangle = 0$, we obtain that

$$\Delta A \, \Delta B \ge \sqrt{\frac{|\langle \psi | [A, B] | \psi \rangle|^2}{4}} = \frac{|\langle \psi | [A, B] | \psi \rangle|}{2}$$

2. This time we will use Cauchy-Schwartz inequality that $|\langle a| b \rangle|^2 \leq \langle a| a \rangle \langle b| b \rangle$ for any vector a and b in any Hilbert space. We also use the following inequality that for any two complex numbers x and y, $|x - y|^2 \leq 2(|x|^2 + |y|^2)$ which one can simply prove.

$$\begin{split} |\langle \psi | [A, B] |\psi \rangle|^{2} &= |\langle \psi | AB |\psi \rangle - \langle \psi | BA |\psi \rangle|^{2} \stackrel{(a)}{\leq} 2(|\langle \psi | AB |\psi \rangle|^{2} + |\langle \psi | BA |\psi \rangle|^{2}) \\ &\stackrel{(b)}{\leq} 2(\langle \psi | AA^{\dagger} |\psi \rangle \langle \psi | B^{\dagger}B |\psi \rangle + \langle \psi | BB^{\dagger} |\psi \rangle \langle \psi | A^{\dagger}A |\psi \rangle) \\ &\stackrel{(c)}{=} 4 \langle \psi | A^{2} |\psi \rangle \langle \psi | B^{2} |\psi \rangle = 4(\Delta A)^{2}(\Delta B)^{2}, \end{split}$$

where (a) follows from the inequality for complex numbers just mentioned, (b) follows by applying the Cauchy-Schwartz inequality and (c) follows from the Hermitian property of A and B.

3. For $\psi = \frac{1}{\sqrt{2}}(|\uparrow\rangle + i |\downarrow\rangle)$ and A = X and B = Z, we have

$$\bar{A} = \langle \psi | A | \psi \rangle = \frac{1}{2} (1 - i) \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 \\ i \end{pmatrix} = 0,$$

$$\bar{A}^2 = \langle \psi | A^2 | \psi \rangle = \frac{1}{2} (1 - i) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 \\ i \end{pmatrix} = 1,$$

$$\bar{B} = \langle \psi | B | \psi \rangle = \frac{1}{2} (1 - i) \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 1 \\ i \end{pmatrix} = 0,$$

$$\bar{B}^2 = \langle \psi | B^2 | \psi \rangle = \frac{1}{2} (1 - i) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 \\ i \end{pmatrix} = 1,$$

where for an operator C, \overline{C} denotes the average of the operator in state $|\psi\rangle$. Therefore, we get

$$(\Delta A)^2 = \langle \psi | A^2 | \psi \rangle - (\langle \psi | A | \psi \rangle)^2 = 1, (\Delta B)^2 = \langle \psi | B^2 | \psi \rangle - (\langle \psi | B | \psi \rangle)^2 = 1.$$

We also know that the Pauli matrices satisfy the identity $[X, Z] = 2 \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$. Thus, we have

$$\langle \psi | [A, B] | \psi \rangle = 2 \frac{1}{2} \begin{pmatrix} 1 & -i \end{pmatrix} \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 \\ i \end{pmatrix} = -2i,$$

which satisfies the uncertainty principle

$$(\Delta A)^2 (\Delta B)^2 = 1 \ge \frac{|2i|^2}{4} = \frac{|\langle \psi | [A, B] | \psi \rangle|^2}{4}.$$

In particular, in this case the uncertainty inequality turns out to be an equality.

4. In this exercise, the state of a particle is represented by a function of coordinate variable $\psi(x)$. We also know how \hat{x} and \hat{p} operators transform this state function. To find the commutator of \hat{x} and \hat{p} , we take an arbitrary state function $\phi(x)$ and investigate how $[\hat{x}, \hat{p}]$ operates on it. Specifically we have

$$\begin{aligned} (\hat{p}\hat{x})\phi(x) &= \hat{p}(\hat{x}\phi(x)) = \hat{p}(x\phi(x)) = -i\hbar\frac{d}{dx}(x\phi(x)) = -i\hbar(\phi(x) + x\frac{d}{dx}\phi(x)), \\ (\hat{x}\hat{p})\phi(x) &= \hat{x}(-i\hbar\frac{d}{dx}\phi(x)) = -i\hbar x\frac{d}{dx}\phi(x), \end{aligned}$$

which implies that $[\hat{x}, \hat{p}]\phi(x) = i\hbar\phi(x)$. As ϕ was an arbitrary function, it results that $[\hat{x}, \hat{p}] = i\hbar$.